Mathematics and Music: Timbre and Consonance

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Introduction & Terminology

Vibrations in Musical Instruments (Timbre)

Consonance & Dissonance

Introduction

Questions:

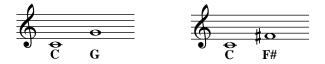
Two musical instruments playing the same note still sound different. Why?

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Introduction

Questions:

- Two musical instruments playing the same note still sound different. Why?
- Some musical intervals sound consonant ("good"?), others dissonant ("bad"?). Why?



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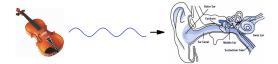


Pressure disturbances in air propagate as waves



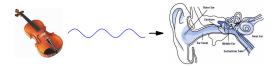
Sound

- Pressure disturbances in air propagate as waves
- Air pressure (within a given frequency band) incident on the ear's basilar membrane is perceived as *sound*



Sound

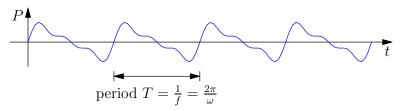
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Music is organized sound

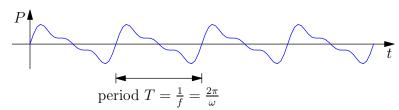
Musical Tones

periodic sound pressure is perceived as a musical tone:



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Musical Tones



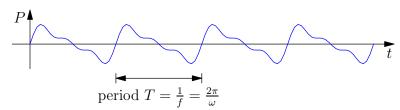
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• the frequency f [Hz = cycles/sec] is perceived as *pitch*

higher $f \Leftrightarrow$ higher pitch

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Musical Tones



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▶ the frequency *f* [Hz = cycles/sec] is perceived as *pitch*

higher $f \Leftrightarrow$ higher pitch

a pure tone of frequency f is sinusoidal:

$$P(t) = A\sin(2\pi ft + \phi)$$

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▶ Pitch perception is logarithmic in frequency:

pitch: $p = \log f$

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Pitch perception is logarithmic in frequency:

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► So translation in pitch is multiplication in frequency:

$$p = p_1 + p_2 \Leftrightarrow f = f_1 f_2$$

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- An pitch interval Δp = p₁ p₂ corresponds to a ratio of frequencies f₁ : f₂.
- A sequence of equally spaced pitches (musical scale)

$$\{p_0, p_0 + \Delta p, p_0 + 2\Delta p, p_0 + 3\Delta p, \ldots\}$$

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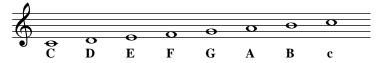
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is a *geometric* sequence in frequency:

$$\{f_0, \alpha f_0, \alpha^2 f_0, \alpha^3 f_0, \ldots\}$$

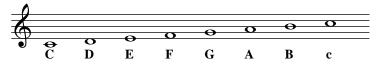
C major scale (equal temperament)



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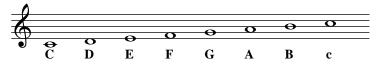
Subset of the 12-tone chromatic scale:

$$f_n = 440 \cdot 2^{n/12}$$
 $(n = -5, -4, \dots, 2)$

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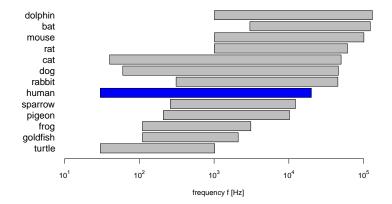
pitch	frequency <i>f_n</i> [Hz]
С	261.6
D	293.6
Е	329.6
F	349.2
G	392.0
Α	440.0
В	493.8
с	523.2

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Limits of Hearing

Hearing Ranges in Various Species



Source: R.Fay, Hearing in Vertebrates: A Psychoacoustics Databook. Hill-Fay Associates, 1988.

 Sound waves are made by a vibrating body (plucked string, hammered block, etc).

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- Sound waves are made by a vibrating body (plucked string, hammered block, etc).
- A freely vibrating body is described by the partial differential equation (*wave equation*)

$$\frac{\partial^2 u}{\partial t^2} = c^2 \nabla^2 u \quad \text{(with initial \& boundary conditions)}$$

 $u(\mathbf{x}, t) =$ displacement from rest at time t at point $\mathbf{x} \in \mathbb{R}^n$.

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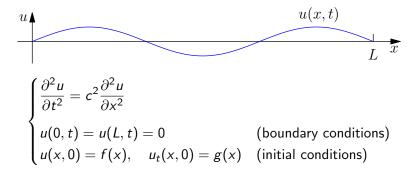
$$\nabla^2 = \text{Laplacian operator} = \begin{cases} \frac{\partial^2 u}{\partial x^2} & \text{on } \mathbb{R} \\\\ \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} & \text{on } \mathbb{R}^2 \end{cases}$$

Example: string fixed at both ends.



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$$\frac{\partial^2 u}{\partial t^2} = c^2 \nabla^2 u$$
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Assume $u(\mathbf{x}, t) = X(\mathbf{x})T(t)$:

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$$= XT'' - c^2 T \nabla^2 X$$

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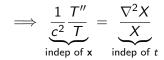
$$\implies \frac{1}{c^2} \frac{T''}{T} = \frac{\nabla^2 X}{X}$$

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$$\implies \underbrace{\frac{1}{c^2} \frac{T''}{T}}_{\text{indep of } \mathbf{x}} = \underbrace{\frac{\nabla^2 X}{X}}_{\text{indep of } t} = \lambda \quad \text{(constant)}$$

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$$\frac{1}{c^2}\frac{T''}{T} = \frac{\nabla^2 X}{X} = \lambda$$

$$T'' + c^2 \lambda T = 0$$

$$\implies T(t) = A \sin(\sqrt{\lambda}ct + \phi) \quad (A, \phi \in \mathbb{R})$$

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 $\nabla^2 X = \lambda X \quad (+ \text{ boundary conditions})$ $\implies \lambda \text{ is an eigenvalue of } \nabla^2$

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By linearity:

$$u(x,t) = \sum_{n=1}^{\infty} A_n \sin(\sqrt{\lambda_n} ct + \phi_n) f_n(x)$$

where λ_n are eigenvalues, f_n are eigenfunctions of ∇^2 .

Summary: motion of a freely vibrating body is

$$u(x,t) = \sum_{n=1}^{\infty} A_n \sin(\underbrace{\sqrt{\lambda_n}c}_{\omega_n} t + \phi_n) f_n(x).$$

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Key points:

Superposition of vibrational modes (pure tones)

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Vibrations in Musical Instruments

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 λ_n are eigenvalues of ∇^2 (for given domain & bc's)

Amplitudes A_n determined by *initial conditions*

Vibrations in Musical Instruments

Summary: motion of a freely vibrating body is

$$u(x,t) = \sum_{n=1}^{\infty} A_n \sin(\underbrace{\sqrt{\lambda_n}c}_{\omega_n} t + \phi_n) f_n(x).$$

Key points:

- Superposition of vibrational modes (pure tones)
- Frequencies are

$$\omega_n = \sqrt{\lambda_n} c$$

 λ_n are eigenvalues of ∇^2 (for given domain & bc's)

- Amplitudes A_n determined by initial conditions
- Smallest λ_n gives the fundamental tone; other modes give upper partials



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Eigenvalue problem $\nabla^2 f = \lambda f$ in 1-D becomes

$$\frac{d^2f}{dx^2} = \lambda f \implies f(x) = A\sin(\sqrt{\lambda}x) + B\cos(\sqrt{\lambda}x).$$

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Boundary condition f(0) = 0 implies B = 0 so

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Boundary condition f(L) = 0 gives

$$0 = A\sin(\sqrt{\lambda}L) \implies \sqrt{\lambda_n}L = n\pi \quad (n = 0, 1, 2, ...)$$

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Frequencies of vibrational modes are given by

$$\lambda_n = \frac{n^2 \pi^2}{L^2} \implies \omega_n = \sqrt{\lambda_n} c = \frac{n \pi c}{L} \quad (n = 1, 2, 3, \ldots)$$

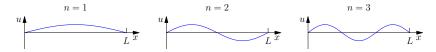
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The eigenfunctions (modes) themselves look like

$$f_n(x)=\sin\frac{n\pi x}{L}.$$



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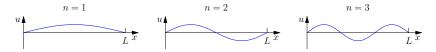
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The string's motion is a superposition of these:

$$u(x,t) = \sum_{n=1}^{\infty} A_n \sin(\omega_n t + \phi_n) f_n(x)$$

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What we hear is a superposition of pure tones:

$$P(t) = \sum_{n=1}^{\infty} A_n \sin(\omega_n t + \phi_n)$$

at discrete frequencies

$$\omega_n = \frac{n\pi c}{L} = n\omega_1 \quad (n = 1, 2, 3, \ldots)$$

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(the harmonic series).

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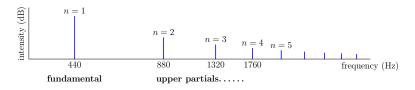
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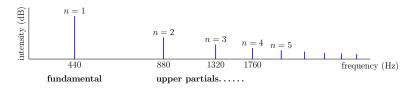
- Frequencies ω_n are integer multiples of the fundamental ω_1 .
- Sound perception is independent of the phase ϕ_n .

For a guitar string playing the note A (440 Hz) we hear:



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For a guitar string playing the note A (440 Hz) we hear:



First 7 upper partials for low C:



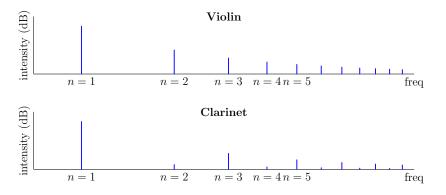
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timbre: the quality or tone distinguishing voices or instruments

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timbre: the quality or tone distinguishing voices or instruments

Timbres of different instruments are distinguished primarily by the frequencies and amplitudes of their spectra:



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Eigenvalues of ∇^2 for various instruments:

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 or $(2n+1)^2$ (depending on bc's)

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 $\lambda_{mn} = n^{\text{th}}$ root of $J_m(\lambda)$, the Bessel function of order m

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Vibrating bars: (e.g. xylophone, marimba)

$$\begin{cases} \lambda_n = (2n+1)^4 & \text{(transverse virbations)} \\ \lambda_m = m^2 & \text{(longitudinal vibrations)} \end{cases}$$

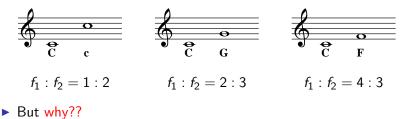
► A musical interval is a given difference in pitch (hence frequency ratio f₁ : f₂) between two tones.

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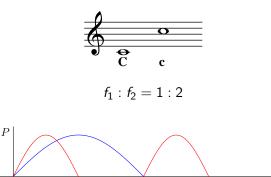
► A musical **interval** is a given difference in pitch (hence frequency ratio f₁ : f₂) between two tones.

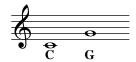
Some intervals sound consonant ("good"?), others dissonant ("bad"?)

- ► A musical interval is a given difference in pitch (hence frequency ratio f₁ : f₂) between two tones.
- Some intervals sound consonant ("good"?), others dissonant ("bad"?)
- Pythagoras: an interval is consonant if the frequencies are in a simple integer ratio:

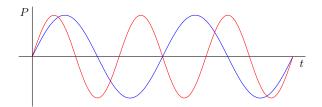


A wrong explanation (Galileo and many others):





 $f_1: f_2 = 3:2$



"The pulses delivered by the two tones ... shall be commensurable in number, so as not to keep the ear-drum in perpetual torment..."

> Galileo Galilei Dialogues Concerning Two New Sciences (1638)

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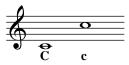
However... for pure tones a mis-tuned interval isn't dissonant!

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- ▶ However... for *pure tones* a mis-tuned interval isn't dissonant!
- ► The reality: dissonance comes from upper partials.

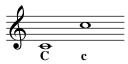
A better explanation:



$$f_1: f_2 = 1:2$$

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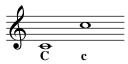
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Consider spectra of upper partials for these tones:

*f*₁: 220 Hz, 440 Hz, 660 Hz, 880 Hz, 1100 Hz, 1320 Hz, ... *f*₂: 440 Hz, 880 Hz, 1320 Hz, 1760 Hz, ...

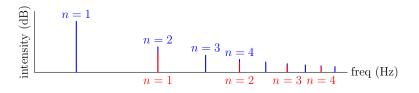
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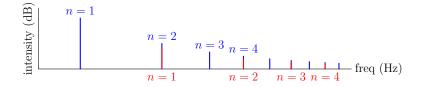


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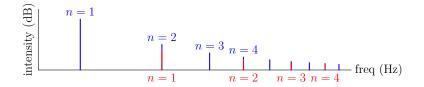
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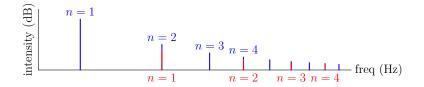
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Upper partials coincide and reinforce each other.



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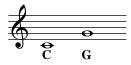
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- The effect is one of altered timbre.



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- Upper partials coincide and reinforce each other.
- The effect is one of altered timbre.
- Invidividual tones are difficult to distinguish.

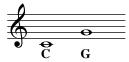
Similarly for a perfect fifth:



 $f_1: f_2 = 3:2$

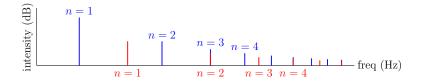
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Similarly for a perfect fifth:



 $f_1: f_2 = 3:2$

Again, some of the upper partials coincide:



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at common multiples of the fundamentals.

So generally. . .

If the fundamentals are in the ratio

$$f_1:f_2=m:n$$

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So generally...

If the fundamentals are in the ratio

$$f_1:f_2=m:n$$

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upper partials coincide for every common multiple of m, n.

Lowest common multiple is mn. The n'th partial of f₁ coincides with the m'th partial of f₂.

So generally...

If the fundamentals are in the ratio

$$f_1:f_2=m:n$$

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- Effect is more audible if the product *mn* is smaller.

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- Simple integer ratios emerge as intervals with strongest mutual reinforement.

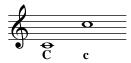
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- Lowest common multiple is mn. The n'th partial of f₁ coincides with the m'th partial of f₂.
- Effect is more audible if the product *mn* is smaller.
- Simple integer ratios emerge as intervals with strongest mutual reinforement.
- But this doesn't really explain dissonance of other intervals.

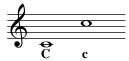
An even better explanation:



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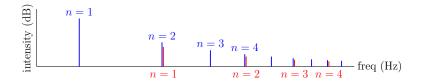


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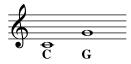
Consider spectra for a slightly mistuned octave:



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Previously coincident partials now differ.

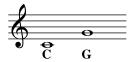
Similarly for the perfect fifth:



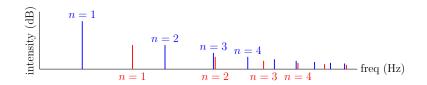
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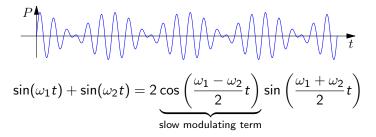
If two pure tones of nearly equal frequency are sounded simultaneously, **beats** are heard:

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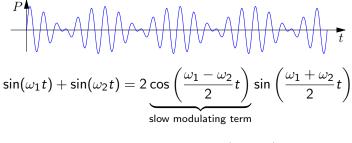
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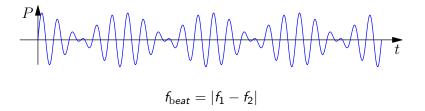
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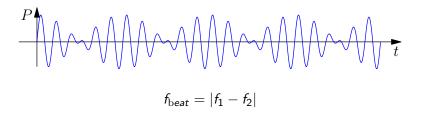
beat frequency:
$$\mathit{f}_{ ext{beat}} = |\mathit{f}_1 - \mathit{f}_2|$$

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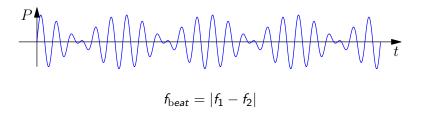


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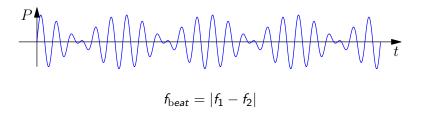


• $f_{\text{beat}} \lesssim 10 \,\text{Hz} \implies$ slow modulation (tremolo), not unpleasant

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*f*_{beat} ≤ 10 Hz ⇒ slow modulation (tremolo), not unpleasant
*f*_{beat} ≥ 50 Hz ⇒ beat frequency becomes an audible tone



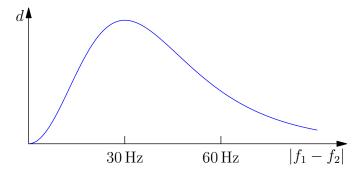
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- $f_{\rm beat}\gtrsim$ 50 Hz \implies beat frequency becomes an audible tone
- $\blacktriangleright~10\,{\rm Hz} \lesssim {\it f}_{{\rm beat}} \lesssim 50\,{\rm Hz}$ gives a "rough", unpleasant feeling

• Maximum dissonance occurs for $|f_1 - f_2| \approx 30 \text{ Hz}$

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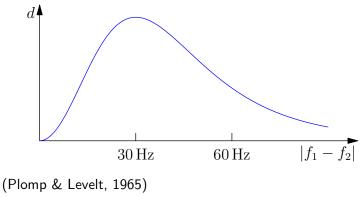
- Maximum dissonance occurs for $|f_1 f_2| \approx 30 \text{ Hz}$
- Dissonance *d* depends (subjectively) on $|f_1 f_2|$:



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(Plomp & Levelt, 1965)

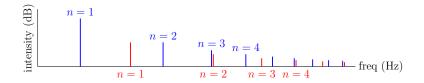
- Maximum dissonance occurs for $|f_1 f_2| \approx 30 \text{ Hz}$
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We can model this with:

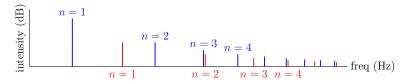
$$d(x) = \frac{(x/30)^2}{(1 + \frac{1}{3}(x/30)^2)^4}$$

When two notes are sounded, dissonance potentially arises from beating between every pair of partials.



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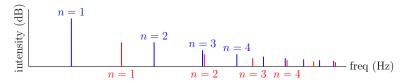


▶ Mistuned fifth (f₁ : f₂ = 3 : 2) with f₁ = 220 Hz and f₂ = 335 Hz:

*f*₁: 220 Hz, 440 Hz, 660 Hz, 880 Hz, 1100 Hz, 1320 Hz, ... *f*₂: 335 Hz, 670 Hz, 1005 Hz, 1340 Hz, 1675 Hz, ...

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 - *f*₁: 220 Hz, 440 Hz, 660 Hz, 880 Hz, 1100 Hz, 1320 Hz, ...
 - *f*₂: 335 Hz, 670 Hz, 1005 Hz, 1340 Hz, 1675 Hz, ...
- Near-concidence of upper partials causes beating:

670 - 660 = 10 Hz and 1340 - 1320 = 20 Hz.

A simple model for relative consonance of intervals:

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A simple model for relative consonance of intervals:

Two notes with fundamental frequencies f₁ and f₂, hence two sequences of partials

$$\{f_1, 2f_1, 3f_1, \ldots\} = \{mf_1 : m = 1, 2, \ldots\}$$

$$\{f_2, 2f_2, 3f_2, \ldots\} = \{nf_2 : n = 1, 2, \ldots\}$$

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$$\{f_2, 2f_2, 3f_2, \ldots\} = \{nf_2 : n = 1, 2, \ldots\}$$

Sum dissonances over all pairs of partials:

total dissonance
$$=\sum_m \sum_n d(|mf_1 - nf_2|)$$
 $d(x) = rac{(x/30)^2}{(1+rac{1}{3}(x/30)^2)^4}$

Fix f_1 and calculate dissonance as a function of f_2 :

total dissonance =
$$\sum_{m} \sum_{n} \underbrace{d(|mf_1 - nf_2|)}_{\text{dissonance of pair }m, n}$$

Summing over the first 7 partials:



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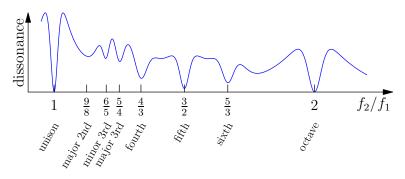
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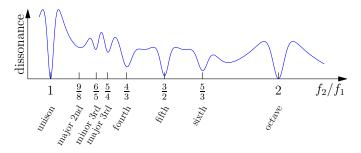
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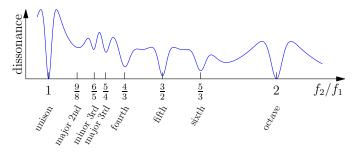
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Consonant intervals are local minima of dissonance.

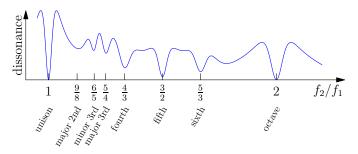


- Consonant intervals are local minima of dissonance.
- Including more partials introduces more local minima.

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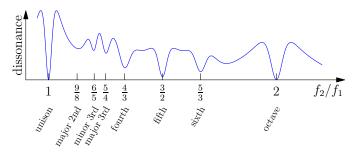


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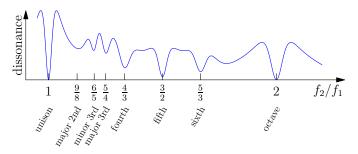
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Strongest consonances are lowest minima.



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- Depth of each minimum determines how well characterized each consonance is (i.e. relative to adjacent intervals).



- Consonant intervals are local minima of dissonance.
- Including more partials introduces more local minima.
- Strongest consonances are lowest minima.
- Depth of each minimum determines how well characterized each consonance is (i.e. relative to adjacent intervals).
- Dissonance curve is a consequence of the underlying timbre (spectrum) of the instrument.